# Calculation of ${ }^{22} \mathrm{Na}(\mathrm{n}, \boldsymbol{\alpha}){ }^{19} \mathrm{~F}$ Cross Sections from the ${ }^{19} \mathrm{~F}(\alpha, \mathrm{n}){ }^{22} \mathrm{Na}$ Cross Sections by Using the Principle of Detailed Balance 

Ali Hassan Ahmed

Dept. of Physics, College of Science, Salahaddin University, Erbil, Iraq
(Received 13/4/2008, Accepted 8/3/2009)


#### Abstract

The cross sections of ${ }^{22} \mathrm{Na}(\mathrm{n}, \alpha)^{19} \mathrm{~F}$ reaction were measured through the ${ }^{19} \mathrm{~F}(\alpha, \mathrm{n}){ }^{22} \mathrm{Na}$ cross sections available in the literature using the principle of detailed balance in the energy range of $\alpha$-particles ( $0.9-6.65$ ) MeV . This work covers for the first time a wide energy range, and a considerable correspondence observed within the resonances of both direct and inverse reactions.


(keywords: ${ }^{22} \mathrm{Na}(\mathrm{n}, \alpha){ }^{19} \mathrm{~F}$, Cross Sections, Principle of Detailed Balance)

## 1. Introduction

Neutron-induced helium-production cross sections play a significant role in the astrophysical processes of nucleosynthesis of light and medium mass nuclei, and it is important for reactor design technology, in particular, for estimating radiation damage, nuclear heating, and induced activity in structural materials. The investigation of the ( $\mathrm{n}, \alpha$ ) reactions is of considerable interest to study reaction mechanisms and nuclear structure, and also they are needed in medical and biological fields to estimate biological effects of neutrons. The inverse of the ${ }^{19} \mathrm{~F}(\alpha, \mathrm{n}){ }^{22} \mathrm{Na}$ reaction may be an important destruction mechanism for ${ }^{22} \mathrm{Na}$ in neutron-rich, high-temperature or explosive nucleosynthesis [1]. At energies between threshold to 10 MeV , the cross section data of the ${ }^{22} \mathrm{Na}(\mathrm{n}, \alpha){ }^{19} \mathrm{~F}$ reaction is rather scanty and are not available due to the instability of ${ }^{22} \mathrm{Na}$. So, in this paper the latest cross sections of ${ }^{19} \mathrm{~F}(\alpha, \mathrm{n})^{22} \mathrm{Na}$ reaction were interpolated and recalculated in fine steps, from which the cross sections of ${ }^{22} \mathrm{Na}(\mathrm{n}, \alpha){ }^{19} \mathrm{~F}$ were calculated by means of the principle of detailed balance, covering the mentioned energy range.

## 2. Theory

A. Cross Section Of ${ }^{19} f(A, N)^{22}$ na

The data of cross section declared by reference [2] have been recalculated in steps of 50 keV up to 10 MeV of the incident $\alpha$-particle energy. As well as, the plot of neutron yields (in units of $n / \alpha$ ) against $\alpha$-energy measured by Nakasima et.al. [3] were used to calculate the cross sections in similar fine steps by means of the formula [4]:
$\sigma(E)=\left[\frac{Y\left(E_{2}\right)-Y\left(E_{1}\right)}{E_{2}-E_{1}}\right] \frac{d E}{d x}$
where $E$ is the mean of the two close energies $E_{1}$ and $E_{2}$, and $(\mathrm{dE} / \mathrm{dx})$ is the stopping power of $\alpha$-particles given by Norman et.al. [5].
The measured cross sections (with their percentage of error) from both references are tabulated in Table (1).

## B. Cross Section Of ${ }^{\mathbf{2 2} \mathrm{na}(\mathrm{N}, \mathrm{A}){ }^{19} \mathrm{f} \text { Reaction: }}$

The principle of detailed balance allows us to calculate the cross section of a reaction from its inverse. The relation between the cross section of the reaction $1+2 \longrightarrow 3+4$ and its inverse is given by $[6,7]:$

$$
\begin{equation*}
\frac{\sigma_{34 \rightarrow 12}}{\sigma_{12 \rightarrow 34}}=\frac{\left(1+\delta_{34}\right) g_{1} g_{2} P_{12}}{\left(1+\delta_{12}\right) g_{3} g_{4} P_{34}} \tag{2}
\end{equation*}
$$

where $\delta_{34}, \delta_{12}$ are the kronecker delta, which accounts for the possibility that the two particles in the exit and entrance channels might be initiated, $g_{i}$ is the spin multiplicity of the ith particle, and $\left(g_{i}=2 I_{i}+1\right.$, where $I_{i}$ is the spin of the ith particle), and $\mathrm{P}_{\mathrm{i}}$ is the momentum of the ith particle, $\quad P_{i}^{2}=2 \mathrm{~m}_{\mathrm{i}} \mathrm{T}_{\mathrm{i}}$

Applying this principle to our reactions,
${ }^{19} \mathrm{~F}+\alpha \longrightarrow{ }^{22} \mathrm{Na}+\mathrm{n} \quad$ direct reaction
${ }^{22} \mathrm{Na}+\mathrm{n} \longrightarrow{ }^{19} \mathrm{~F}+\alpha$ inverse reaction we obtain:

$$
\begin{array}{r}
\begin{array}{r}
\frac{\sigma_{n, \alpha}}{\sigma_{\alpha, n}}=\frac{P_{\alpha}^{2}}{P_{n}^{2}} \frac{\left(2 I_{\alpha}+1\right)\left(2 I_{19}{ }_{F}+1\right)}{\left(2 I_{n}+1\right)\left(2 I_{22}+1\right)}--- \\
\\
=\frac{m_{\alpha} E_{\alpha}}{m_{n} E_{n}} \frac{(2 \times 0+1)\left(2 \times \frac{1}{2}+1\right)}{\left(2 \times \frac{1}{2}+1\right)(2 \times 3+1)} \\
\text { en } \quad \sigma_{n, \alpha}=0.56689 \frac{E_{\alpha}}{E_{n}} \sigma_{\alpha, n}
\end{array}
\end{array}
$$

Then
(4)

The value of $E_{n}$ in the inverse reaction ${ }^{22} \mathrm{Na}(\mathrm{n}, \alpha){ }^{19} \mathrm{~F}$ corresponding to $\mathrm{E}_{\alpha}$ in the direct reaction is given approximately by [8]:

$$
E_{n}(l a b .) \cong \frac{M\left({ }^{19} F\right)}{\left.M^{(22} N a\right)}\left(E_{\alpha}-E_{t h r}\right)
$$

(5)
where $\mathrm{E}_{\mathrm{thr}}=2.36 \mathrm{MeV}$ for the ${ }^{19} \mathrm{~F}(\alpha, \mathrm{n})^{22} \mathrm{Na}$ reaction, Using eqs. (4) \& (5) we obtain the cross sections of the inverse reaction ${ }^{22} \mathrm{Na}(\mathrm{n}, \alpha)^{19} \mathrm{~F}$ as tabulated in Table (2).

## 3. Results And Discussion:

The reproduced cross sections of ${ }^{19} \mathrm{~F}(\alpha, \mathrm{n})^{22} \mathrm{Na}$ reaction in the energy range ( $2.4-10 \mathrm{MeV}$ ) are listed in Table (1). The corresponding cross sections of ${ }^{22} \mathrm{Na}(\mathrm{n}, \alpha)^{19} \mathrm{~F}$ reaction calculated using the principle of detailed balance in the energy range $(0.9-6.65 \mathrm{MeV})$ are listed in Table (2) and plotted in Figure (1). The lower energy regions, near threshold, were not presented due to the disagreement between the cross section data of the ${ }^{19} \mathrm{~F}(\alpha, \mathrm{n}){ }^{22} \mathrm{Na}$ reaction for the two references in this energy range. These results and especially in this energy range were not done previously except the work of Wrean [1] in the energy range $(0-700 \mathrm{keV})$; so there is no available data for comparison.

Through Tables (1 and 2) and Figure (1), the detailed correspondence between the reaction ${ }^{19} \mathrm{~F}(\alpha, \mathrm{n}){ }^{22} \mathrm{Na}$ and its inverse ${ }^{22} \mathrm{Na}(\mathrm{n}, \alpha){ }^{19} \mathrm{~F}$ are observed revealing a good approve of our calculations using the reciprocity theorem.

## 5. Conclusions:

The cross section measurement of nuclear reactions for the unstable target nuclei which are incapable experimentally can be performed using the principle of detailed balance. This method reveals its high adequacy in the calculation of ${ }^{22} \mathrm{Na}(\mathrm{n}, \alpha)^{19} \mathrm{~F}$ cross section from the ${ }^{19} \mathrm{~F}(\alpha, \mathrm{n}){ }^{22} \mathrm{Na}$ reaction cross sections.

Table (2) : The adopted Cross Section of ${ }^{22} \mathrm{Na}(\mathrm{n}, \square){ }^{19} \mathrm{~F}$ Reaction Calculated from the ${ }^{19} \mathrm{~F}(\square, \mathrm{n}){ }^{22} \mathrm{Na}$ Reaction.

| $\mathbf{E}_{\mathbf{n} \square}(\mathbf{M e V})$ | ( $\square_{\mathbf{n}, \square}$ in mb) |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | Nakasima et.al. 1992 [3] (10\%) | Shibata and Matsunobu 2003 [2] (10\%) | Weighted Average $\pm$ Error |  |
| 0.9 | $54.96100 \pm 5.496100$ | $48.08090 \pm 4.808090$ | 51.0636 | $\pm 3.6188$ |
| 0.95 | $57.33550 \pm 5.733550$ | $54.49060 \pm 5.449060$ | 55.8407 | $\pm 3.9498$ |
| 1 | $59.11440 \pm 5.911440$ | $57.97890 \pm 5.797890$ | 58.5356 | $\pm 4.1393$ |
| 1.05 | $61.20040 \pm 6.120040$ | $64.38110 \pm 6.438110$ | 62.7102 | $\pm 4.4357$ |
| 1.1 | $62.96310 \pm 6.296310$ | $68.23800 \pm 6.823800$ | 65.3888 | $\pm 4.6274$ |
| 1.15 | $66.05170 \pm 6.605170$ | $72.64660 \pm 7.264660$ | 69.0363 | $\pm 4.8871$ |
| 1.2 | $69.98470 \pm 6.998470$ | $78.41030 \pm 7.841030$ | 73.7206 | $\pm 5.2212$ |
| 1.25 | $74.65630 \pm 7.465630$ | $83.58690 \pm 8.358690$ | 78.6192 | $\pm 5.5681$ |
| 1.3 | $77.61680 \pm 7.761680$ | $88.13650 \pm 8.813650$ | 82.2117 | $\pm 5.8249$ |
| 1.35 | $81.65120 \pm 8.165120$ | $93.32690 \pm 9.332690$ | 86.7134 | $\pm 6.1452$ |
| 1.4 | $85.78260 \pm 8.578260$ | $99.39290 \pm 9.939290$ | 91.5928 | $\pm 6.4941$ |
| 1.45 | $90.36180 \pm 9.036180$ | $103.2189 \pm 10.32189$ | 95.9402 | $\pm 6.7989$ |
| 1.5 | $94.66890 \pm 9.466890$ | $106.9966 \pm 10.69966$ | 100.082 | $\pm 7.0901$ |
| 1.55 | $99.99630 \pm 9.999630$ | $110.7512 \pm 11.07512$ | 104.8263 | $\pm 7.422$ |
| 1.6 | $105.5954 \pm 10.55954$ | $114.5469 \pm 11.45469$ | 109.7078 | $\pm 7.7639$ |
| 1.65 | $108.2471 \pm 10.82471$ | $117.9528 \pm 11.79528$ | 112.6843 | $\pm 7.9753$ |
| 1.7 | $113.9749 \pm 11.39749$ | $123.4455 \pm 12.34455$ | 118.333 | $\pm 8.3741$ |
| 1.75 | $119.9108 \pm 11.99108$ | $127.4576 \pm 12.74576$ | 123.4542 | $\pm 8.7336$ |
| 1.8 | $123.7133 \pm 12.37133$ | $132.1228 \pm 13.21228$ | 127.6419 | $\pm 9.0305$ |
| 1.85 | $130.6263 \pm 13.06263$ | $134.1826 \pm 13.41862$ | 132.3567 | $\pm 9.3599$ |
| 1.9 | $136.1019 \pm 13.61019$ | $138.0203 \pm 13.80203$ | 137.0477 | $\pm 9.691$ |
| 1.95 | $141.3770 \pm 14.13770$ | $139.7445 \pm 13.97445$ | 140.5513 | $\pm 9.9386$ |
| 2 | $146.9303 \pm 14.69303$ | $140.9251 \pm 14.09251$ | 143.8025 | $\pm 10.1706$ |
| 2.05 | $152.3665 \pm 15.23665$ | $145.7371 \pm 14.57371$ | 148.9045 | $\pm 10.5317$ |
| 2.1 | $158.4026 \pm 15.84026$ | $153.9572 \pm 15.39572$ | 156.1167 | $\pm 11.0402$ |
| 2.15 | $164.8331 \pm 16.48331$ | $160.6854 \pm 16.06854$ | 162.7085 | $\pm 11.5062$ |
| 2.2 | $170.1389 \pm 17.01389$ | $168.1487 \pm 16.81487$ | 169.1321 | $\pm 11.9597$ |
| 2.25 | $175.1453 \pm 17.51453$ | $178.5950 \pm 17.85950$ | 176.8269 | $\pm 12.5041$ |
| 2.3 | $180.9686 \pm 18.09686$ | $187.3928 \pm 18.73928$ | 184.0687 | $\pm 13.0176$ |
| 2.35 | $185.7008 \pm 18.57008$ | $193.4562 \pm 19.34562$ | 189.4199 | $\pm 13.3968$ |
| 2.4 | $190.5797 \pm 19.05797$ | $197.8783 \pm 19.78783$ | 194.0919 | $\pm 13.7268$ |
| 2.45 | $198.0207 \pm 19.80207$ | $204.5616 \pm 20.45616$ | 201.1849 | $\pm 14.2278$ |
| 2.5 | $203.4998 \pm 20.34998$ | $212.4128 \pm 21.24128$ | 207.7654 | $\pm 14.6946$ |
| 2.55 | $208.2430 \pm 20.82430$ | $218.8230 \pm 21.88230$ | 213.271 | $\pm 15.0852$ |
| 2.6 | $213.9576 \pm 21.39576$ | $223.8816 \pm 22.38816$ | 218.6948 | $\pm 15.468$ |
| 2.65 | $219.1259 \pm 21.91259$ | $229.9099 \pm 22.99099$ | 224.2591 | $\pm 15.8621$ |
| 2.7 | $223.7978 \pm 22.37978$ | $238.0469 \pm 23.80469$ | 230.4831 | $\pm 16.3054$ |
| 2.75 | $228.0188 \pm 22.80188$ | $243.7030 \pm 24.37030$ | 235.34 | $\pm 16.6502$ |
| 2.8 | $234.3834 \pm 23.43834$ | $245.9491 \pm 24.59491$ | 239.8879 | $\pm 16.9676$ |

Table (2) Continued

| $\mathbf{E}_{\mathbf{n} \square}(\mathrm{MeV})$ | ( $\square_{\mathrm{n}, \square}$ in mb) |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | Nakasima et.al. 1992 [3] (10\%) | Shibata and Matsunobu 2003 [2] (10\%) | Weighted Average $\pm$ Error |  |
| 4.85 | $332.6140 \pm 33.26140$ | $343.7158 \pm 34.37158$ | 337.9827 | $\pm 23.9022$ |
| 4.9 | $333.4375 \pm 33.34375$ | $344.1036 \pm 34.41036$ | 338.6027 | $\pm 23.9458$ |
| 4.95 | $334.1212 \pm 33.41212$ | $345.6168 \pm 34.56168$ | 339.6746 | $\pm 24.0221$ |
| 5 | $334.9638 \pm 33.49638$ | $347.2314 \pm 34.72314$ | 340.8771 | $\pm 24.1075$ |
| 5.05 | $335.6622 \pm 33.56622$ | $349.4705 \pm 34.94705$ | 342.2882 | $\pm 24.2083$ |
| 5.1 | $335.6879 \pm 33.46879$ | $351.3875 \pm 35.13875$ | 343.1791 | $\pm 24.2728$ |
| 5.15 | $336.1857 \pm 33.61857$ | $353.6666 \pm 35.36666$ | 344.4834 | $\pm 24.3665$ |
| 5.2 | $337.0611 \pm 33.70611$ | $355.6876 \pm 35.56876$ | 345.8739 | $\pm 24.4658$ |
| 5.25 | $337.3776 \pm 33.73776$ | $356.5656 \pm 35.65656$ | 346.4414 | $\pm 24.5065$ |
| 5.3 | $338.1412 \pm 33.81412$ | $359.3990 \pm 35.93990$ | 348.1229 | $\pm 24.6274$ |
| 5.35 | $338.7646 \pm 33.87646$ | $358.8387 \pm 35.88387$ | 348.2245 | $\pm 24.6334$ |
| 5.4 | $339.2732 \pm 33.92732$ | $360.9924 \pm 36.09924$ | 349.4598 | $\pm 24.7224$ |
| 5.45 | $339.5859 \pm 33.95859$ | $361.6681 \pm 36.16681$ | 349.9323 | $\pm 24.7562$ |
| 5.5 | $339.5193 \pm 33.95193$ | $361.9644 \pm 36.19644$ | 350.0244 | $\pm 24.7631$ |
| 5.55 | $339.3821 \pm 33.93821$ | $363.7999 \pm 36.37999$ | 350.7441 | $\pm 24.8163$ |
| 5.6 | $339.9098 \pm 33.99098$ | $366.6802 \pm 36.66802$ | 352.2822 | $\pm 24.928$ |
| 5.65 | $340.8289 \pm 34.08289$ | $373.9598 \pm 37.39589$ | 355.862 | $\pm 25.1903$ |
| 5.7 | $341.3892 \pm 34.13892$ | $377.5791 \pm 37.75791$ | 357.6671 | $\pm 25.3229$ |
| 5.75 | $342.2945 \pm 34.22945$ | $378.7547 \pm 37.87547$ | 358.6857 | $\pm 25.3953$ |
| 5.8 | $342.7775 \pm 34.27775$ | $380.5566 \pm 38.05566$ | 359.6992 | $\pm 25.4692$ |
| 5.85 | $341.9173 \pm 34.19173$ | $381.9815 \pm 38.19815$ | 359.7388 | $\pm 25.4763$ |
| 5.9 | $342.6084 \pm 34.26084$ | $382.4201 \pm 38.24201$ | 360.3347 | $\pm 25.5179$ |
| 5.95 | $343.8505 \pm 34.38505$ | $383.0640 \pm 38.30640$ | 361.348 | $\pm 25.5883$ |
| 6 | $342.5726 \pm 34.25726$ | $384.5978 \pm 38.45978$ | 361.1646 | $\pm 25.5808$ |
| 6.05 | $342.9837 \pm 34.29837$ | $384.1378 \pm 38.41378$ | 361.2389 | $\pm 25.5843$ |
| 6.1 | $345.4842 \pm 34.54842$ | $385.7220 \pm 38.57220$ | 363.3955 | $\pm 25.7348$ |
| 6.15 | $345.2757 \pm 34.52757$ | $386.2159 \pm 38.62159$ | 363.4616 | $\pm 25.7408$ |
| 6.2 | $346.3491 \pm 34.63491$ | $388.2819 \pm 38.82819$ | 364.9297 | $\pm 25.8464$ |
| 6.25 | $345.4777 \pm 34.54777$ | $389.9077 \pm 38.99077$ | 365.0182 | $\pm 25.8577$ |
| 6.3 | $343.6108 \pm 34.36108$ | $391.4934 \pm 39.14934$ | 364.4464 | $\pm 25.8249$ |
| 6.35 | $343.9536 \pm 34.39536$ | $393.7367 \pm 39.37367$ | 365.5007 | $\pm 25.9036$ |
| 6.4 | $343.1968 \pm 34.31968$ | $394.4645 \pm 39.44645$ | 365.2847 | $\pm 25.8918$ |
| 6.45 | $342.3385 \pm 34.23385$ | $395.6925 \pm 39.56925$ | 365.1785 | $\pm 25.8894$ |
| 6.5 | $341.5059 \pm 34.15059$ | $396.7368 \pm 39.67368$ | 365.0123 | $\pm 25.8824$ |
| 6.55 | $340.1263 \pm 34.01263$ | $398.6062 \pm 39.86062$ | 364.7657 | $\pm 25.8735$ |
| 6.6 | $339.3204 \pm 33.93204$ | $399.9859 \pm 39.99859$ | 364.7084 | $\pm 25.8755$ |
| 6.65 | $338.1565 \pm 33.81565$ | $401.3959 \pm 40.13959$ | 364.4078 | $\pm 25.8616$ |



Figure (1): Cross Section as a function of neutron energy for the ${ }^{22} \mathrm{Na}(\mathrm{n}, \mathrm{a}){ }^{19} \mathrm{~F}$ Reaction.

## References:

[1] P.R. Wrean, Ph.D. thesis, California Institute of Technology, (1998).
[2] K.Shibata and T.Matsunobu., Japanese Evaluated Nuclear Data Library, JENDL / AN-2003,JAERI, Japan Atomic Energy Research Institute (JAERI), (2003).
[3] R. Nakasima, H. Matsunobu, T. Oku, S. Iijima, Y. Naito, and F. Masukawa, "Data Book for Calculating Neutron Yields from ( $\alpha, n$ ) Reaction and Spontaneous Fission", JAERI 1324 (1992) (in Japanese).
[4] K.Shibata, JAERI-M 83-221 (1983).
[5] E.B. Norman, T.E.Chupp, K.T.Lesko, P.J.Grant, and Gene L. Woodruff, Physical Review, C30(1984)1375.
[6] W.A. Fowler, G.R. Caughlan, and B.A.Zimmerman, Annu. Rev. Astron. Astrophys. 5 (1967) 525.
[7] E. Segre, Nuclei and Particles, $2^{\text {nd }}$ edition, Benjamin / Cummings Publishing Company, Inc.(U.S.A, 1979), Page 510.
[8] R.L. Macklin and J.H. Gibbons, Physical Review, v.165, No. 4 (1968).

تم حساب الدقاطع العرضية للتفاعل النووي البيانات النووية NDC باستخدام مبدأ النوازن التفصيلي في مدى طاقة جسيمات ألفا (6.65-0.9) MeV). في هذه الدراسة ولأول مرة تم تغطية مدى طاقي واسع للقائف، حيث وجدت توافق جيد بين الرنائن المقاسة من الثناعلات المباثرة والعكسية.

